

Chiral Phase transition

$SU(N_c)$ gauge theory w/ N_f quarks

$$\begin{aligned} \mathcal{L} &\supset \sum_{i=1}^{N_f} \bar{\psi}^i (i\not{\partial}) \psi_i \\ &= \sum_{i=1}^{N_f} \left\{ \bar{\psi}_L^i (i\not{\partial}) \psi_{L,i} + \bar{\psi}_R^i (i\not{\partial}) \psi_{R,i} \right\} \end{aligned}$$

Global Symmetry:

$$\begin{aligned} G_F &= U(N_f)_L \times U(N_f)_R \\ &= U(1)_V \times U(1)_A \times SU(N_f)_L \times SU(N_f)_R \end{aligned}$$

Topic: T -dependence of manifest G_F symmetry?

Comments:

- $U(1)_A$ is anomalous

$$\partial_\mu J^{\mu 5} \sim g^2 N_f \text{tr } \tilde{F} F$$

↑
 $SU(N_c)$ gauge fields

But anomaly still respects $Z_{N_c} \subset U(1)_A$

Why? $\mathcal{Z} \rightarrow e^{i\alpha \gamma_5/2} \mathcal{Z}$ under $U(1)_A$

In instanton background,

$$\underbrace{\langle \mathcal{Z}_1 \dots \mathcal{Z}_n \rangle}_{2N_F} \neq 0$$

$\downarrow U(1)_A$

$$\underbrace{e^{i\alpha N_F}}_{\text{wavy}} \langle \dots \rangle$$

$$= 1 \quad \text{if } \alpha = \frac{2\pi M}{N_F}$$

- Anomaly vanishes in Planar limit:

$$N_c \rightarrow \infty, \quad \Lambda = g^2 N_c \text{ fixed}$$

Because quark loops give subleading
rel. to glue:

$$\partial_\mu T^{\mu 5} \sim \frac{\Lambda}{N_c} \rightarrow 0$$

So, in presence of anomaly:

$$G_F' = U(1)_V \times Z_{N_F} \times SU(N_F)_L \times SU(N_F)_R$$

Aside: Spontaneous breaking of Z_{NF} gives no NG bosons, but in $N_c \rightarrow \infty$ limit, there is one, η' . Instantons give it a mass

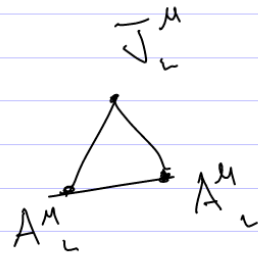
Does T-dependent vacuum respect G_F' ?

Already answered @ $T=0$:

Strong dynamics leads to spontaneously broken chiral symm. in IR
+ 't Hooft anomaly matching:

G_F' has 't Hooft anomalies:

$$\left(\text{SU}(N_F)_{L,R} \right)^3$$



Conserved along RG flow. But, in IR we have confinement. "Can check" that confined spectrum can't reproduce anomaly unless G_F' is spontaneously broken

What about $T > 0$?

Lattice Shows a $T_{ch} < \infty$

↑

: independent of g_2

at which G_F' is restored.

Assume so, try to learn if 1st
or 2nd order

Effective theory of order parameter:

$$\phi^i_j \sim \langle \bar{\psi}_L^i \psi_{R,j} \rangle$$

Transforms under $U(1)_A \times SU(N_f)_L \times SU(N_f)_R$

as

$$\phi \rightarrow e^{i\alpha} U_L^+ \phi U_R$$

Most general, renormalizable $\mathcal{L}(\phi)$:

$$\begin{aligned} \mathcal{L}_\phi = & \frac{1}{2} \text{tr}(\partial_\mu \phi^+ \partial^\mu \phi) - \frac{1}{2} M_\phi^2 \text{tr}(\phi^+ \phi) \\ & - \frac{\pi^2}{3} g_1 (\text{tr} \phi^+ \phi)^2 - \frac{\pi^2}{3} g_2 \text{tr}(\phi^+ \phi)^2 \end{aligned}$$

$$g_2 \geq 0, \quad g_1 + g_2/N_f \geq 0 \quad \text{for stability}$$

Does not reproduce anomaly:

$$\mathcal{L}'_{\phi} = c (\det \phi + \det \phi^{\dagger})$$

$$\sim \underbrace{\langle \bar{\psi}_L \psi_R \dots \bar{\psi}_L \psi_R \rangle}_{2N_f}$$

Finite T , compactify Euclidean time,
 'integrate out all massive (EWT) modes
 and focus on 3D theory

Since phase transition occurs @ $T = T_{ch}$

$$M_{\phi}^2(T) \sim M_{\phi}^2(T - T_{ch})$$

$$\uparrow$$

$$T < T_{ch}, M_{\phi}^2(T) < 0$$

↓
 minimum is at $\phi \neq 0$

SSB

$$T > T_{ch} \quad M_{\phi}^2(T) > 0, \quad \langle \phi \rangle = 0$$

$$@ T = T_{cu}, \quad M_\phi^2 = 0$$

"widely believed that" if an IR stable fixed point exists for $\mathcal{L}_\phi(T=T_{cu})$ then transition is 2nd order. Else, 1st order

Why?

Speculation:
 • 2nd order described by CFT, i.e. fixed point?
 • 2nd order exhibits universality?

Stable fixed point means independent of UV

Anyway: calculate β -functions of g_1, g_2

in $4-\epsilon$ dimensions:

For $C=0$

$$\beta_1 = -\epsilon g_1 + \frac{(N_F^2 + 4)}{3} g_1^2 + \frac{4N_F}{3} g_1 g_2 + g_2^2$$

$$\beta_2 = -\epsilon g_2 + 2g_1 g_2 + \frac{2N_F}{3} g_2^2$$

No stable fp if $N_F > \sqrt{3}$

Stable if $0 < N_F \leq \sqrt{2}$

$$c \neq 0$$

W/o details, for

$N_F \geq 3$ c doesn't affect 1st order

$N_F = 2$, c can push transition to
2nd order

What is c ?

c couples to anomaly of $U(1)_A \sim$ instantons

Assume $c(\tau) \sim d_{\frac{1}{2}}(\tau)$



T -dependent instanton density

$c(0) \sim \Theta(1)$ since $c(0)$

gives mass to η^i , separating it from
other NGs.

$(CT \rightarrow \infty) \rightarrow 0$ since instantons of size

$\rho > \frac{1}{T}$ don't fit into
compact time

So, at high T , approx.
restoration of $U(1)_A$.

Since anomaly gives η' mass apart from
other NGs, this has effects on spectrum